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# Revealing the Role of  $TiO<sub>2</sub>$  Surface Treatment of Hematite Nanorods Photoanodes for Solar Water Splitting

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**S** Supporting Information

[AB](#page-5-0)STRACT: Ultrathin  $TiO<sub>2</sub>$  is deposited on conventional hydrothermal grown hematite nanorod arrays by atomic layer deposition (ALD). Significant photoelectrochemical water oxidation performance improvement is observed when the ALD TiO<sub>2</sub>-treated samples are annealed at 650 °C or higher temperatures. The electrochemical impedance spectroscopy (EIS) study shows a surface trap-mediated charge transfer process exists at the hematite−electrolyte interface. Thus, one possible reason for the improvement could be the increased surface states at the hematite surface, which leads to better



charge separation, less electron−hole recombination, and hence, greater improvement of photocurrent. Our Raman study shows the increase in surface defects on the ALD TiO<sub>2</sub>-coated hematite sample after being annealed at 650  $^{\circ}$ C or higher temperatures. A photocurrent of 1.9 mA cm<sup>−</sup><sup>2</sup> at 1.23 V (vs RHE) with a maximum of 2.5 mA cm<sup>−</sup><sup>2</sup> at 1.8 V (vs RHE) in 1 M NaOH under AM 1.5 simulated solar illumination is achieved in optimized deposition and annealing conditions.

KEYWORDS: photoelectrochemical cells (PEC), hematite, atomic layer deposition (ALD), nanorods, TiO<sub>2</sub>, surface treatment

## 1. INTRODUCTION

Photoelectrochemical cells (PEC) offer the ability to harness solar photons to split water, thereby producing hydrogen as a chemical fuel.<sup>1−3</sup> Hematite ( $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>) is a very promising material for PEC water splitting because of its favorable optical band gap (∼2[.1](#page-5-0) [e](#page-6-0)V), excellent stability under caustic operating conditions, earth abundance, and low cost.<sup>4,5</sup> The theoretical solar-to-hydrogen efficiency based on a semiconductor with this band gap can achieve a value of  $16.8\%$ . H[ow](#page-6-0)ever, the overall water splitting efficiency of hematite photoanodes falls well short of the theoretical maximum effic[ie](#page-6-0)ncy, primarily due to the very short electron−hole pair lifetime (<10 ps) and hole diffusion length (2−4 nm), which gives rise to a high recombination rate of photogenerated carriers in the bulk.<sup>4</sup> Considerable efforts have been devoted to improve the actual PEC performance of hematite by employin[g](#page-6-0) nanostructuring strategies, $7-15$  as well as introducing dopant atoms into hematite.16−<sup>25</sup> In addition, the kinetics of the interfacial extraction [of](#page-6-0) holes from the hematite surface for water oxidatio[n \(oxy](#page-6-0)gen evolution) reaction seems to be sluggish, which results in increased charge recombination with a concomitant loss of photocurrent and hence the efficiency. In most of the reports, the surface traps have been suggested as an additional loss mechanism in various hematite photoanodes.26−<sup>29</sup> Indeed, improved hematite-based photochemical water splitting has been reported through passivating of surface sta[tes u](#page-6-0)sing very thin passivating layers, such as  $\text{Fe}_x\text{Sn}_{1-x}\text{O}_4$ <sup>30</sup>ALD  $\text{Al}_2\text{O}_3$ <sup>26</sup> and ALD  $\text{TiO}_2$ <sup>26,31</sup> However, a

surface trap-mediated charge transfer process at the hematiteelectrolyte interface has been suggested by Bisquert and their co-workers based on their electrochemical impedance spectroscopy (EIS) study, in which the water oxidation takes place predominantly from surface trapped holes, not directly from valence band holes.<sup>32</sup> In such a charge transfer mode, the surface states help the charge injection at the hematite and electrolyte interface[.](#page-6-0) Thus, there could be a new surface treatment strategy to improve the photoelectrochemical water oxidation by increasing those surface states and consequently enhancing the charge transfer at the hematite/electrolyte interface.

In this work, we report a novel surface treatment strategy to improve the charge transfer of hematite photoanode by coating the hematite surface with an ultrathin  $TiO<sub>2</sub>$  layer using ALD and followed by an appropriate heat treatment at 650 °C. We perform a systematic investigation of the role of the treatments by accurately controlling the thickness and deposition conditions of the  $TiO<sub>2</sub>$  overlayer. During the annealing process, as the solid-state diffusion reaction occurs between the hematite and the  $TiO<sub>2</sub>$  overlayer, Ti is gradually incorporated into hematite from the surface. Simultaneously, small amount of Fe is diffused into the  $TiO<sub>2</sub>$  layer leading to the amorphization of TiO<sub>2</sub>. The increased Raman intensity of the LO mode of

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<span id="page-1-0"></span>hematite shows the increase of surface defects at hematite surface after the solid reaction. EIS measurements show the presence of a trap-mediated charge transfer at the hematiteelectrolyte, which can be shifted to higher potentials with the ALD  $TiO<sub>2</sub>$  layer. With optimized ALD  $TiO<sub>2</sub>$  thickness and post annealing time at 650 °C, a photocurrent of 1.9 mA cm<sup>−</sup><sup>2</sup> at 1.23 V (vs RHE) is achieved. This novel surface treatment strategy can be further applied to a broader range of nanostructured hematite photoanodes prepared by other methods

## 2. EXPERIMENTAL SECTION

2.1. Preparation of Hematite Nanorod Arrays and ALD  $TiO<sub>2</sub>$ Coating. Hematite nanorod arrays were grown on fluorine-doped tin oxide substrates (FTO, thickness 2.2 mm, resistance <14 ohm/square, transmittance >90%, purchased from Wuhan Geao, China) using a hydrothermal method. Typically, a 50 mL aqueous solution containing 10 mmol of FeCl<sub>3</sub>·6H<sub>2</sub>O (Sigma, ACS reagent 97%) and 15 mmol of urea (Sigma, 98%) was transferred into a 100 mL glass bottle and 12 substrates  $(1 \text{ cm} \times 2.5 \text{ cm})$  were placed vertically standing at the bottom of the bottle, then sealed with a cap. After 10 h of reaction at 100 °C, a uniform layer of yellow FeOOH film was formed on the FTO substrate. The substrates were thoroughly rinsed in distilled water (DI) water and blow-dried with  $N_2$  gun. After annealing at 550  $^{\circ}$ C under air for 2 h, FeOOH was transformed to hematite. ALD TiO<sub>2</sub> was conducted on a Fiji 200 system with Tetrakis dimethylamino titanium (TDMATi) and  $H_2O$  as precursors at 120 °C. The TDMATi source was kept at 70 °C to achieve a suitable vapor pressure. The precursors were pulsed in the reactor with 40 sccm constant Ar flow, and the base pressure for deposition was 0.8 Torr. The pulse time for Ti and  $H_2O$  were 0.6 and 0.1 s, respectively. The purge time was 10 s for both precursors. The growth rate per cycle of  $TiO<sub>2</sub>$  was 0.06 nm, which is inconsistent with previous report.<sup>33</sup> After ALD TiO<sub>2</sub> coating, the hematite samples were kept horizontally inside a tube furnace in ambient condition and heated up to 650 °[C](#page-6-0) with a heating rate of 25  $\mathrm{^{\circ}C/min}$ , and kept at this temperature for different times (e.g., 30 min, 60 min, 90 min).

2.2. Characterization of Samples. Morphology of samples was characterized by a JEOL, JSM-7600F field emission scanning electron microscope (FESEM), high-resolution transmission electron microscopy (HRTEM) and scanning TEM (JEOL 2100F). The hematite nanorods were scraped from the FTO substrate and followed by a ultrasonic treatment in ethanol for TEM sample preparation. Raman spectra were collected on Renishaw inVia Raman system with 532 nm excitation wavelength, and a 50× objective lens was used in our measurement. Diffuse reflectance ultraviolet−visible (UV−vis) absorption spectra of the films were obtained using an Agilent Cary series UV−vis−NIR spectrophotometer (model Cary 500) equipped with an external diffuse reflectance accessory (DRA-2500). The diameter of integrating sphere is 150 mm. X-ray photoemission spectroscopy (XPS) was collected on a home-built ultra high vacuum (UHV) system with a monochromatic X-ray source at  $hv = 1486.7 \text{ eV}$ (Al,  $K\alpha$ ) from SPECS and Omircon electron analyzer (EA125).

2.3. Photoelectrochemical Property Measurements. PEC measurements were carried out using CHI 660D workstation (CH Instruments, Inc.) with a three-electrode electrochemical system in 1 M NaOH (pH 13.6) electrolyte. Platinum mesh and Ag/AgCl were employed as counter and reference electrodes, respectively. The working surface area was 0.3 cm<sup>2</sup>. The light source was simulated sunlight from a 150 W xenon solar simulator (67005, Newport Corp.) through a solar filter with a measured intensity equivalent to standard AM1.5 sunlight (100 mW cm<sup>−</sup><sup>2</sup> ) on the sample surface. IPCE characteristics were measured with a xenon light source (MAX-302, Asahi Spectra Co. Ltd.) coupled with a monochromator (CMS-100, Asahi Spectra Co. Ltd.) from 305 to 630 nm, at a potential of 1.23 V vs RHE. A Si photodiode (Bentham, DH-Si) with known IPCE was used to calculate the IPCE of the pure hematite and  $TiO<sub>2</sub>$ -treated hematite photoanodes. A source meter (Keithley Instruments Inc., Model:

2400) was used to measure the photocurrent of Si diode. The working station (CHI 660D, CH Instruments, Inc.) mentioned above was used to measure the photocurrent of samples. The EIS measurements were carried out using an Autolab PGSTAT 302F in a three-electrode electrochemical system. A sinusoidal voltage perturbation of 20 mV was applied from 100 kHz to 0.1 Hz, with a DC voltage sweep over the range of interest. The same reference and counter electrodes used for PEC measurements were used.

## 3. RESULTS AND DISCUSSION

Figure 1a, b shows the top-view and cross-section SEM images of the as grown hematite nanorod arrays coated with 40 cycles



Figure 1. SEM image of the ALD  $TiO<sub>2</sub>$ -coated hematite after being annealed at  $650 \degree$ C for 30 min: (a) top view and (b) cross-section image. (c) HRTEM of 40 cycles of ALD  $TiO<sub>2</sub>$ -coated hematite after annealed at 650 °C for 30 min. (d) Schematic effect of ultrathin ALD  $TiO<sub>2</sub>$  overlayer on hematite nanorods for efficient solar water splitting.

of ALD TiO<sub>2</sub> overlayer and annealed at 650  $^{\circ}$ C for 30 min. As the TiO<sub>2</sub> over layer is very thin, ~2.5 nm for 40 cycles, morphology remains similar to the pristine hematite nanorods arrays. The morphology of the hydrothermal FeOOH and the pristine hematite without ALD  $TiO<sub>2</sub>$  coating are shown in Figure S1. The diameter of the nanorods is around 60 to 80 nm, and the length is ∼700 nm. The aspect ratio of the [hematite](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf) nanorods is around 10. Figure 1c shows a typical HRTEM image of the hematite nanorod arrays sample with 40 ALD TiO<sub>2</sub> cycles and post annealed at 650  $^{\circ}$ C for 30 min. The lattice of the hematite core can be clearly seen. The amorphous overlayer around the hematite nanorod surface could also be observed. The thickness of the overlayer is around 2 nm, which is consistent with the estimated ALD  $TiO<sub>2</sub>$  thickness (2.4 nm for 40 ALD cycles). To further investigate the ultrathin amorphous overlayer, Raman, XRD, and STEM measurements were carried out. Figure 2 shows the Raman spectra of the  $TiO<sub>2</sub>$ coated hematite samples annealed at different temperatures (400 and 650  $^{\circ}$ [C\) and t](#page-2-0)he pristine hematite nanorod arrays annealed at 650 °C for comparison. All the Raman peaks of the pristine sample: 224, 244, 291, 297, 410, 499, 610, and 660  $cm^{-1}$  can be assigned to hematite.<sup>34</sup> After ALD TiO<sub>2</sub> and further annealing at 400 °C for 30 min, a small but obvious Raman peak at 143 cm<sup>−</sup><sup>1</sup> (red marke[d](#page-6-0) in Figure 2) assigned to anatase TiO<sub>2</sub> E<sub>g</sub> mode is observed,<sup>35</sup> and all the Raman peaks assigned to hematite remained with no [obvious](#page-2-0) change. This suggests that the nanorod structu[re](#page-6-0) consist of hematite core

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Figure 2. Raman spectra of the pristine hematite nanorod arrays (black curve), ALD TiO<sub>2</sub>-coated hematite annealed at 400 °C (red curve), and 650 °C (blue curve), respectively.

and  $TiO<sub>2</sub>$  shell after the 400 °C annealing. By increasing the post annealing temperature to 650 °C, the Raman peak of TiO<sub>2</sub> at ∼143 cm<sup>−</sup><sup>1</sup> disappears and the intensity of the hematite LO mode peak (∼660 cm<sup>−</sup><sup>1</sup> ) dramatically increases (Figure 2). It should be noted that the LO mode of hematite at 660  $cm^{-1}$  is normally forbidden by the selection rules, so the origin of this band is attributed to the breakdown of Raman symmetry selection rules induced by short-range order or lattice defects, as in the case of nanodimensions due to a large amount of surface defects. $36-38$  All these observations indicate occurrence <sup>-38</sup> All these observations indicate occurrence of solid reaction between  $TiO<sub>2</sub>$  overlayer and hematite core when annealed [at 65](#page-6-0)0 °C. During the annealing process a small amount of Fe ions diffused into the  $TiO<sub>2</sub>$  layer leading to amorphization of  $TiO<sub>2</sub>$  layer. However, this cannot be observed by XRD due to the overlap of the XRD peaks of  $TiO<sub>2</sub>$  and SnO2 (from FTO substrate, Figure S2). Simultaneously, some Ti atoms on the surface diffused into the hematite core,  $14$  which may increase the surface [defects an](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf)d lattice distortion of hematite. The XRD pattern of the hematite nanorods [sh](#page-6-0)ows a preferential orientation of the (110) peak, as the hematite growth direction is along (110) direction, as seen the HRTEM image in Figure 1c. The absence of peak shift or broadening of hematite in Raman and XRD (Figure 2 and Figure S2) studies also impl[ies that t](#page-1-0)he modifications due to  $ALD TiO<sub>2</sub>$  treatment are limited to the surface of the hematite [nanorods](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf). STEM images and elemental mapping of the hematite nanorod also show a uniform distribution of Ti (see Figure S3).

XPS was used to verify the Ti incorporation of the ALD  $TiO<sub>2</sub>$ coated hematite nanorods arrays. Figu[re 3a show](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf)s the survey scan of the TiO<sub>2</sub> coated hematite after annealed at 650 °C together with the bare hematite. Both spectra shows the same features of the hematite in Fe 2p and the existence of O. But only the ALD  $TiO<sub>2</sub>$  coated sample shows the Ti signal. The binding energies of Ti  $2p_{1/2}$  and  $2p_{3/2}$  are observed at 464.5 and 458.3 eV, respectively, indicating a normal state of  $Ti^{4+.25}$  The Fe  $2p_{3/2}$  peak of the ALD TiO<sub>2</sub> treated hematite sample slightly shift to the lower energy side compared to that for t[he](#page-6-0) bare hematite sample, as seen in Figure 3b. This is consistent with expected binding energy shift based on the fact that electronegativity of Ti (1.54) is lower than that of Fe (1.83). The substitution of Fe with Ti in O−Fe−O will thereby lower the binding energy. Figure S4 shows the O 1s peak of the bare  $Fe<sub>2</sub>O<sub>3</sub>$  and the ALD TiO<sub>2</sub> treated Fe<sub>2</sub>O<sub>3</sub>. The O 1s peak for the bare Fe<sub>2</sub>O<sub>3</sub> sample [is center](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf)ed around 529.9 eV, which is



Figure 3. (a)XPS survey of the bare hematite nanorod arrays after being annealed at 650 °C and the ALD TiO<sub>2</sub> 40 cycle coated hematite sample annealed at 650 °C. (b) Fe 2p details for the bare hematite and ALD TiO<sub>2</sub>-coated hematite sample after being annealed at 650 °C. (c) Ti 2p chemical details for the ALD-coated hematite.

consistent with the reported value.<sup>25</sup> After ALD  $TiO<sub>2</sub>$  treatment the O 1s peak is shifted to slightly higher energy, which is consistent with pervious study [of](#page-6-0) the Ti doped hematite.<sup>25</sup> From the XPS spectra, there is no Sn signal detected, which indicates the Sn diffusion from FTO substrate effect [is](#page-6-0) minimum (and perhaps nonexistent) in our study. Diffused reflectance UV−vis spectra and corresponding Tauc-Plots of our samples are shown in Figure S5. The band gap obtained from Tauc-Plots for all the samples are almost the same, ∼2.05 eV. The slight change of ba[nd gap imp](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf)lies that Ti doping in the hematite nanorods is insignificant.

Figure 4 shows the photocurrent−potential (J−V) curves of the pristine hematite and hematite coated with 40 cycles ALD TiO<sub>2</sub> layers annealed at different temperatures. Negligible p[hotocurre](#page-3-0)nt is observed for the pristine hematite sample annealed at 650 °C, which is consistent with the previous report.<sup>24</sup> The photoresponse of the hematite coated with ALD  $TiO<sub>2</sub>$  and annealed at 400 °C is similar to the one of the pristin[e h](#page-6-0)ematite. The only difference is the potential at which the current rise is shifted from ∼1.5 to ∼1.7 V vs RHE, similar to the observation of the dark current (Figure S6). This could be due to the formation of the heterojunction between hematite and  $TiO<sub>2</sub>$ , as  $TiO<sub>2</sub>$  has much [lower val](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf)ence band,<sup>39</sup> which blocks the hole transfers and hence higher bias is needed for dark current onset. When the annealing temperature [is](#page-6-0) increased to 650 °C, significant photocurrent can be observed. With this annealing temperature a solid reaction occurs between the hematite and  $TiO<sub>2</sub>$  over layer. When the annealing temperature is further increased to 750 °C only a little change in J−V curve is observed. Thus, the occurrence of the solid

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Figure 4. Photocurrent vs potential curves of 40 ALD  $TiO<sub>2</sub>$  cycle coated hematite photoanodes with different post-annealing temperatures.

reaction is more critical to the PEC improvement than the unintentional Sn doping from FTO, which usually occurs after annealed at 750  $^{\circ}$ C.<sup>30</sup>

To investigate into the role of the ALD TiO<sub>2</sub> surface treatment on the P[EC](#page-6-0) performance, we systematically carried out a variation of  $TiO<sub>2</sub>$  thickness and annealing time. PEC was measured on hematite samples with different ALD  $TiO<sub>2</sub>$ thicknesses (from 20 to 160 cycles) and annealing times (30 to 90 min) at 650 °C, the results are shown in Figure 5. For 30 min annealing treatment at 650 °C, the photocurrent density improves when the ALD  $TiO<sub>2</sub>$  cycles are increased from 20 to 80. However, a further increase in the ALD  $TiO<sub>2</sub>$  cycles beyond 80 results in a reduction of the photocurrent density (see Figure 5a, c). The current density for the 80 cycles ALD  $TiO<sub>2</sub>$ coated sample reaches ∼1.1 mA cm<sup>−</sup><sup>2</sup> at 1.23 V vs RHE in 1 M NaOH electrolyte (pH 13.6). A photocurrent onset potential of  $\sim$ 0.9 V vs RHE is observed for the samples with 20 ALD TiO<sub>2</sub> cycle numbers; while for the samples with thicker  $TiO<sub>2</sub>$ overlayers the onset potential shifted positively, reaching ∼1.1 V vs RHE for the sample with 160 ALD cycles  $TiO<sub>2</sub>$  overlayer as shown in Figure 5a. The shifting of the onset potential with thicker ALD  $TiO<sub>2</sub>$  overlayer could be due to the formation of a heterojunction between  $Fe<sub>2</sub>O<sub>3</sub>$  and TiO<sub>2</sub>, which eventually inhibits the charge transport across the junction and limits the overall PEC performance. Figure 5b shows the annealing time effect with the ALD cycles fixed to 40 cycles. The photocurrent improves when the annealing time increases from 30 to 60 min, and then decreases slightly for 90 min annealing time. The increase of the photocurrent when the annealing process is prolonged to 60 min could be due to the further increase of the defects states as the solid state reaction at the  $Fe<sub>2</sub>O<sub>3</sub>/TiO<sub>2</sub>$ progresses, which will be discussed later. When the annealing time is longer than 60 min, the slight decrease in the photocurrent may be a result of the loss in the conductivity of FTO substrate.<sup>40</sup> Figure 5d shows the photocurrent densities under key potentials versus the annealing time. At the optimized ALD [Ti](#page-6-0)O<sub>2</sub> coating of 40 cycles ( $\sim$ 2.4 nm) and annealing time of 60 min at 650  $^{\circ}$ C, the photocurrent at 1.23 V vs RHE is ∼1.9 mA cm<sup>−</sup><sup>2</sup> , which is close to the highest reported photocurrent of 2.4 mA cm<sup>−</sup><sup>2</sup> at 1.23 V for hematite prepared by APCVD without using any catalyst, $12$  and reaches a plateau of about 2.4 mA cm<sup>−</sup><sup>2</sup> at 1.4−1.7 V before the current increases exponentially, see Figure 5b. IPCE [me](#page-6-0)asurement shows the optimized sample can reach 25%, see Figure S7. The I−T curve in Figure S8 shows excellent stability of our hematite photoanode.

[The electrica](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf)l effects of the surface [treatment](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf) with ALD  $TiO<sub>2</sub>$ and post annealing were studied by EIS. The pristine  $Fe<sub>2</sub>O<sub>3</sub>$  and Fe<sub>2</sub>O<sub>3</sub> coated with 40 ALD TiO<sub>2</sub> (post annealed at 650 °C for 1 h) samples were analyzed under illumination with a DC



Figure 5. Photocurrent vs potential curves of hematite photoanodes with (a) different ALD TiO<sub>2</sub> cycles (650 °C annealing time: 30 min fixed), and (b) different annealing time at 650 °C (ALD TiO<sub>2</sub> 40 cycles fixed). (c, d) Current densities for hematite photanodes with different ALD TiO<sub>2</sub> cycles and annealing time at 650 °C, respectively.

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Figure 6. (a) Nyquist plots obtained at 1.15 V Vs RHE under illumination for the pristine F<sub>2</sub>O<sub>3</sub> and the 40 ALD cycle TiO<sub>2</sub>-coated F<sub>2</sub>O<sub>3</sub> with 650  $^{\circ}$ C annealing samples. The lines represent the fitting employing the equivalent circuit in the inset. (b) Total charge transfer resistance ( $R_{\text{trapping}}$  +  $R_{\text{ct, trap}}$ ) and (c)  $C_{\text{trap}}$  and  $R_{\text{ct, trap}}$  for both analyzed samples under illumination. (d) Mott–Schottky analysis obtained from  $C_{\text{bulk}}$  under dark conditions.



Figure 7. Transient photocurrent of the hematite nanorods arrays samples with (a) different ALD TiO<sub>2</sub> cycles and (b) different annealing time at 650 °C, potential 1.0 V vs RHE, 1 M NaOH.

voltage sweep across the working region. Representative Nyquist plots are shown in Figure 6a, featuring two semicircles in the impedance spectrum. The measured spectra were fitted to the equivalent circuit in the inset of Figure 6a comprising a series resistance  $(R_s)$ , a bulk capacitance corresponding the  $Fe<sub>2</sub>O<sub>3</sub>$  (C<sub>bulk</sub>), a trapping resistance accounting for the charge transfer from the bulk to the traps  $(R_{\text{trapping}})$ , the capacitance of the traps  $(C_{trap})$  and the charge transfer from the traps to the electrolyte  $(R_{\text{ct, trap}})$ . The valley of  $R_{\text{ct, trap}}$  at a similar potential than the peak of  $C_{trap}$  (Figure 6c) confirms the presence of a surface states mediated charge transfer, following a previously reported trend with  $Fe<sub>2</sub>O<sub>3</sub>$  systems.<sup>32</sup> In such a charge transfer process, the water oxidation takes place predominantly from surface trapped holes rather than [fr](#page-6-0)om valence band holes. After being annealed at 650 °C, the ALD TiO<sub>2</sub> layer changes the distribution of the surface states, reflected on the  $C_{trap}$ , which reduces both  $R_{\text{trapping}}$  and  $R_{\text{ct, trap}}$ . This results in a better

total charge transfer resistance (Figure 6b) and enhances the performance. Figure 6d shows the Mott−Schottky plots of the pristine  $Fe<sub>2</sub>O<sub>3</sub>$  and the 40 ALD cycles  $TiO<sub>2</sub>$  treated sample. The slopes determined from the Mott−Schottky plots are used to estimate the carrier densities using the equation;  $N_D = (2/$  $e_0 \varepsilon e_0$ )[d(1/C<sup>2</sup>)/dV]<sup>-1</sup>, where  $e_0$  is the electron charge (1.602  $\times$  $10^{-19}$  C),  $\varepsilon$  is the dielectric constant of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> (80),  $\varepsilon_0$  is the permittivity of vacuum (8.854 ×  $10^{-14}$  Fm<sup>-1</sup>), N<sub>D</sub> is the donor density, and  $V$  is the potential applied at the electrode.<sup>30</sup> The as-calculated carrier densities for pristine  $Fe<sub>2</sub>O<sub>3</sub>$  and the 40 ALD cycle TiO<sub>2</sub>-treated sample are similar, 2.13  $\times$  10<sup>[18](#page-6-0)</sup> and 5.49  $\times$  10<sup>18</sup> cm<sup>-3</sup>, respectively. The estimated flat band potential of the pristine  $Fe<sub>2</sub>O<sub>3</sub>$  is around 0.45 V vs RHE, which is in line with previous reported values for nanostructured hematite.<sup>24</sup> No obvious flat band potential shift is observed for the ALD  $TiO<sub>2</sub>$ -treated sample. The similar carrier density and flat ban[d p](#page-6-0)otential indicate that the Ti doping effect in this

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Figure 8. Schematics showing ALD TiO<sub>2</sub> overlayer thickness effect under 650 °C annealing, (a) with thin ALD over layers and (b) with thick ALD  $TiO<sub>2</sub>$  over layers.

study is insignificant. The dominant reason for the increase in photocurrent is the improved charge transfer on the hematite surface via the increased surface states after the ALD  $TiO<sub>2</sub>$ treatment.

To further investigate the ALD  $TiO<sub>2</sub>$  thickness and annealing time effect, we recorded transient photocurrents of the hematite samples under key bias, as these measurements have been reported to substantiate the effect of electron−hole recombination at the hematite surface in previous studies.<sup>26,41</sup> The transient photocurrent at 1.0 V vs RHE of hematite photoanodes with diff[erent](#page-6-0) ALD  $TiO<sub>2</sub>$  cycles and different annealing time at 650 °C are shown in Figure 7a, b, respectively. When light is turned on, photogenerated holes travel to the semiconductor−liquid junction [\(SCLJ\)](#page-4-0) and accumulate because of the slow OER kinetics, $29$  or because of trap states in the bulk and on the surface.<sup>26–29,42</sup> This induces a sharp anodic current spike, which decays un[til](#page-6-0) equilibrium is eventually reached between water o[xidation](#page-6-0) and charge recombination. Conversely, the cathodic transient peak observed when the light is turned off has been assigned to electrons diffusing from the external circuit and recombining with the accumulated holes at the SCLJ. The transient photocurrent spike decreases when the ALD  $TiO<sub>2</sub>$  layer thickness increases at annealing conditions fixed at 650 °C for 30 min (Figure 7a). The smaller transient photocurrent spike means less hole accumulation at the SCLJ. When the ALD TiO<sub>2</sub> is [40 cycle](#page-4-0)s and the annealing time at 650  $^{\circ}$ C is increased from 30 to 60 min, the photocurrent spike becomes much smaller (Figure 7b). When the annealing time is further increased to 90 min, the photocurrent spike looks similar to the 60 min one.

Our experim[ents](#page-4-0) [lead](#page-4-0) us to conclude the effect of ALD  $TiO<sub>2</sub>$ overlayer thickness to the PEC performance of hematite photoanodes annealed at 650 °C, as shown in Figure 8. First, the Ti diffusion into the  $Fe<sub>2</sub>O<sub>3</sub>$  increases the surface defect states, which enhances the trap-mediated charge transfer process at the hematite-electrolyte interface and hence improves the photocurrent. Because thermal diffusion is a concentration and time dependent process, the density of the defect states depends on both the  $TiO<sub>2</sub>$  thickness and the annealing time. This explains why the photocurrent increases with increasing annealing time (Figure 5b) at fixed ALD cycles. Second, if the  $TiO<sub>2</sub>$  layer is too thick, it will block the hole transfer. That could be th[e reason](#page-3-0) for the decreasing photocurrent for the thicker  $TiO<sub>2</sub>$  overlayer samples (Figure 5a), even it shows smaller transient photocurrent spike (Figure

7a). A moderate thickness of ALD  $TiO<sub>2</sub>$  over layer is therefore preferred because it can create sufficient defect states without [b](#page-4-0)locking the hole transfer. Here, the optimized ALD  $TiO<sub>2</sub>$ thickness is about 2.4 nm with 40 ALD cycles annealed at 650 °C for 60 min.

#### 4. CONCLUSIONS

In summary, an effective surface treatment by ALD of hematite nanorods for efficient water oxidation is reported. The increased PEC performance is mainly attributed to the enhanced hole transfer via surface trap at the hematite− electrolyte interface when an amorphous  $TiO<sub>2</sub>$  layer is observed. This enhanced PEC performance is observed when the ALD  $TiO<sub>2</sub>$  layer thickness is below a critical thickness. If the thickness is too thick, the hole transfer from hematite to liquid is blocked because of formation of a substantive heterojunction. A photocurrent of 1.9 mA cm<sup>-2</sup> at 1.23 V (vs RHE) under AM 1.5 simulated solar illumination is achieved in optimized conditions.

## ■ ASSOCIATED CONTENT

#### **6** Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsami.5b01394.

SEM images, XRD, XPS, dark current curves, IPCE [spectra, and](http://pubs.acs.org) I−T curve (PD[F\)](http://pubs.acs.org/doi/abs/10.1021/acsami.5b01394)

#### ■ AUTHOR I[N](http://pubs.acs.org/doi/suppl/10.1021/acsami.5b01394/suppl_file/am5b01394_si_001.pdf)FORMATION

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Notes

The auth[ors declare no competi](mailto:Lydiawong@ntu.edu.sg)ng financial interest.

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